



# EFFECTS OF MOLECULAR WEIGHT AND INTERACTION PARAMETER ON THE GLASS TRANSITION TEMPERATURE OF POLYSTYRENE MIXTURES AND ITS BLENDS WITH POLYSTYRENE/POLY (2,6-DIMETHYL-*p*-PHENYLENE OXIDE)

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**Abstract**—The glass transition temperature ( $T_g$ ) of mixtures of polystyrene (PS) with different molecular weight and of blends of poly(2,6-dimethyl-*p*-phenylene oxide) (PPO) and polystyrene with different molecular weight (DMWPS) was studied by a DSC method. For the whole range of composition, the curves of  $T_g$  vs composition obtained by experiment were compared with predictions from the Fox, Gordon–Taylor, Couchman and Lu–Weiss, equations. It was found that the experimental results were not in agreement with those from the Fox, Gordon–Taylor and Couchman equations for the binary mixtures of DMWPS, where the interaction parameter  $\chi$  was approximately zero. However, for the blends PPO/DMWPS ( $\chi < 0$ ), with an increase of molecular weight of PS, it was shown that the experimental results fitted well with those obtained from the Couchman, Gordon–Taylor and Fox equations, respectively. Furthermore, the Gordon–Taylor equation was nearly identical to the Lu–Weiss equation when  $|\chi|$  was not very large. Further, the dependence of the change of heat capacity associated with the glass transition ( $\Delta C_p$ ) on the molecular weight of PS was investigated and an empirical equation was presented. © 1997 Elsevier Science Ltd

## INTRODUCTION

The glass transition temperature of polymer mixtures can be calculated from values of the pure components. Many empirical formulae suggested by Gordon–Taylor [1], Fox [2], Jenckel–Heusch [3] and Simha–Boyer [4] *et al.* were successively presented and extensively used in practice. Couchman [5] suggested a theoretical equation from classical thermodynamics, which could be reduced approximately to the Gordon–Taylor or Fox equations [5]. However, some of the problems of the thermodynamics in Couchman's derivation were discussed by Goldstein [6]. According to the Gibbs–DiMarzio molecular theory of the glass transition [7], Gordon [8] *et al.* obtained a formula, whose form was identical to that of the Gordon–Taylor equation. By using statistical thermodynamics, they provided a physical meaning of the adjustable parameter in the Gordon–Taylor equation. These formed the theoretical foundation of the Gordon–Taylor equation. Taking the specific interaction between polymer components into account, Kwei [9] introduced an interaction term into the Gordon–Taylor equation

and obtained a generalized Kwei equation. By adjusting two parameters of this equation, this generalized equation [10] can fit well with all curves of the glass transition temperature ( $T_g$ ) vs composition, absent of a cusp, for polymer blends [11]. Somewhat later, Lu and Weiss [12, 13] presented a specific formula by combining the Flory–Huggins interaction parameter ( $\chi$ ) [14] with the Kwei equation. Painter *et al.* [15] and Braun *et al.* [16] also

Table 1. List of the monodisperse polystyrene studied

| PS samples | $M_w (\times 10^{-4})$ | $\eta$ | $T_g$ (K) | $\Delta C_p$ (J/g K) |
|------------|------------------------|--------|-----------|----------------------|
| PS-1       | 0.0859                 | 1.15   | 282       | 0.466                |
| PS-2       | 0.29                   | 1.05   | 335       | 0.406                |
| PS-3       | 0.36                   | 1.06   | 345       | 0.396                |
| PS-4       | 0.40                   | 1.04   | 343       | 0.396                |
| PS-5       | 0.51                   | 1.04   | 358       | 0.389                |
| PS-6       | 0.56                   | 1.03   | 357       | 0.387                |
| PS-7       | 0.71                   | 1.03   | 360       | 0.385                |
| PS-8       | 1.00                   | 1.03   | 358       | 0.381                |
| PS-9       | 2.09                   | 1.03   | 366       | 0.370                |
| PS-10      | 2.24                   | 1.03   | 374       | 0.367                |
| PS-11      | 12.9                   | 1.06   | 379       | 0.359                |
| PS-12      | 16.8                   | 1.02   | 380       | 0.355                |
| PS-13      | 34.1                   | 1.07   | 380.5     | 0.352                |
| PS-14      | 58.2                   | 1.11   | 381       | 0.352                |

$T_g$  and  $\Delta C_p$  are the average values of three experimental results.  $\eta$  is the polydispersity index.

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Table 2. Glass transition data of the binary mixtures of PS of different molecular weights

|            | 0/100 | 20/80 | 40/60 | 50/50 | $\Delta C_{p1}$<br>60/40 | $\Delta C_{p2}$<br>80/20 | 100/0 | (J/g K) | (J/g K) | $k = \frac{\Delta C_{p1}}{\Delta C_{p2}}$ |
|------------|-------|-------|-------|-------|--------------------------|--------------------------|-------|---------|---------|---|
| PS-1/PS-11 | 379   | 350   | 325   | 311.5 | 289                      | 282                      | 0.359 | 0.466   |         | 1.30                                      |
| PS-2/PS-11 | 379   | 365   | 356   | 349   | 341                      | 335                      | 0.359 | 0.406   |         | 1.13                                      |

$\Delta C_{p1}$  is the change in heat capacity associated with glass transition for PS-11.  
 $\Delta C_{p2}$  is the change in heat capacity associated with glass transition for PS-1 or PS-2.

Table 3. Glass transition data of the blends of PPO with different molecular weight PS

|           | 0/100 | 20/80 | 40/60 | 50/50 | $\Delta C_{p1}$<br>60/40 | $\Delta C_{p2}$<br>80/20 | 100/0 | (J/g K) | (J/g K) | $k = \frac{\Delta C_{p1}}{\Delta C_{p2}}$ |
|-----------|-------|-------|-------|-------|--------------------------|--------------------------|-------|---------|---------|---|
| PS-1/PPO  | 459   | 417   | 369   | 349   | 330                      | 302                      | 284   | 0.247   | 0.456   | 1.85                                      |
| PS-5/PPO  | 489   | 451   | 425   | 411   | 396                      | 375                      | 361   | 0.247   | 0.392   | 1.59                                      |
| PS-10/PPO | 489   | 457   | 432   | 420   | 410                      | 391                      | 376   | 0.247   | 0.364   | 1.47                                      |
| PS-11/PPO | 489   | 461   | 434   | 425   | 413                      | 396                      | 379   | 0.247   | 0.359   | 1.45                                      |
| PS-14/PPO | 489   | 466   | 438   | 425   | 417                      | 399                      | 382   | 0.247   | 0.352   | 1.43                                      |

$\Delta C_{p1}$  is the change in heat capacity associated with glass transition for PPO.  
 $\Delta C_{p2}$  is the change in heat capacity associated with glass transition for the different molecular weight PS.

performed much significant work in this respect. However, the dependence of these equations (for example, the Gordon-Taylor, Fox and Couchman equations *et al.*) on the interaction parameter and molecular weight is seldom studied. In this paper, the effects of the interaction parameter and molecular weight on the relationship of  $T_g$  and composition are presented, and the adaptability of the empirical and theoretical equations are also discussed.

EXPERIMENTAL

The molecular parameters of monodisperse PS used in this work are listed in Table 1. Each PS sample was dissolved in benzene to form a solution with a concentration of 5 mg/mL. Blends PS-1/PS-11 and PS-2/PS-11 with the weight ratio of 0/100, 20/80, 40/60, 50/50, 60/40, 80/20 and

100/0 were prepared from benzene solution (5 mg/mL). A similar method was used to prepare blends of poly(2,6-dimethyl-*p*-phenylene oxide) ( $M_n = 4.83 \times 10^4$ ,  $M_w/M_n = 1.723$ ) and PS with different molecular weights (PS-1, PS-5, PS-10, PS-11 and PS-14) from chloroform solution (5 mg/mL). The solvent was removed in air by evaporation. The cast films were dried in a vacuum oven at  $85 \pm 2^\circ\text{C}$  for 3 days and then at a temperature  $20\text{--}30^\circ\text{C}$  higher than  $T_g$  of the respective sample for 36 hr.

Thermal analysis of the samples was performed with a Perkin-Elmer DSC-2C calorimeter with a heating rate of  $20^\circ\text{C}/\text{min}$  by twice scanning, and the sample weight was 10 mg. The samples were first scanned until the temperature was  $30^\circ\text{C}$  higher than the  $T_g$ , then quickly quenched to  $-40^\circ\text{C}$  and scanned for the second time.  $T_g$  was defined as

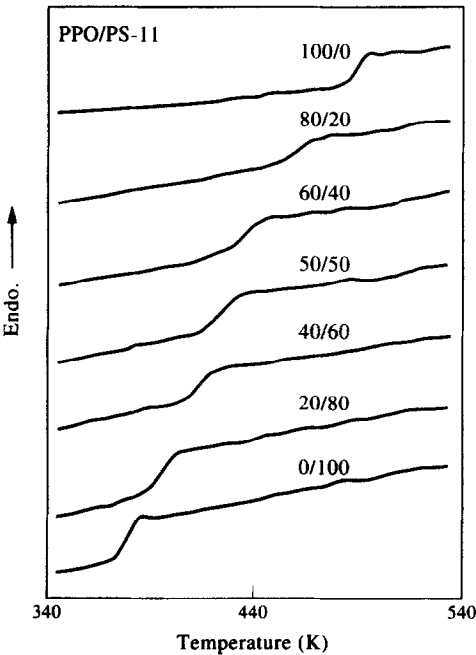


Fig. 1. DSC scans for PPO/PS-11 blends.

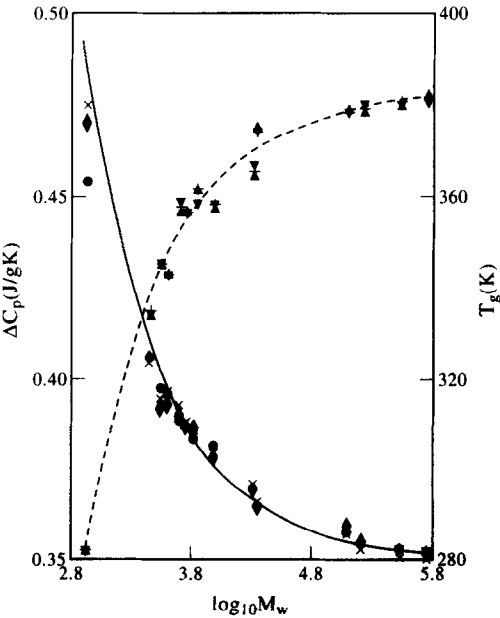


Fig. 2.  $\Delta C_p$  and  $T_g$  vs  $\log_{10} M_w$  for DMWPS:  $\blacktriangle$ ,  $\blacktriangledown$  and  $+$  represent the first, second and third experimental values of  $T_g$ , respectively.  $\bullet$ ,  $\blacklozenge$  and  $\times$  represent the first, second and third experimental values of  $\Delta C_p$ , respectively. --- the fitted curve of the average  $T_g$  vs  $\log_{10} M_w$ . — the calculated curve from equation (1).

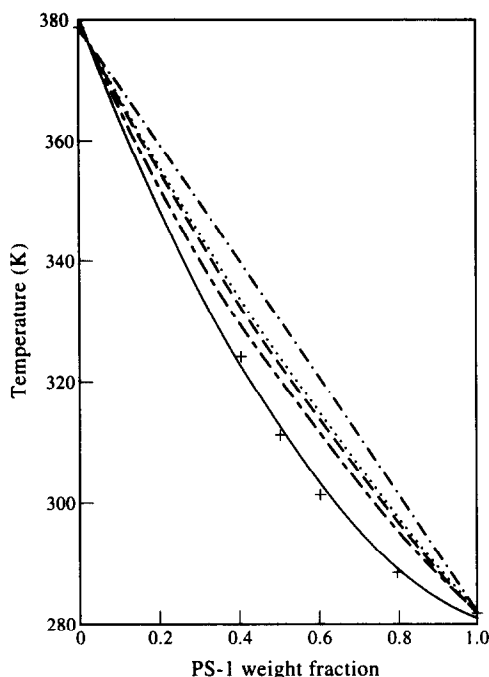


Fig. 3. Dependence of  $T_g$  on composition in PS-1/PS-11 mixtures: + experimental values; --- additive curve; — experimental curve; ··· Fox equation curve; -·- Gordon-Taylor equation curve for  $k = 1.30$ ; - - - Couchman equation curve for  $k = 1.30$ .

the midpoint of the change in specific heat for the second scanning and  $\Delta C_p$  was the change in heat capacity associated with the glass transition with values given in Tables 1, 2 and 3. The DSC thermograms are shown in Fig. 1 (with only the DSC diagram of blends PPO/PS-11 cited as an example).

## RESULTS AND DISCUSSION

### Dependence of $\Delta C_p$ on molecular weight

The experimental values of  $\Delta C_p$  for different molecular weights are shown in Fig. 2.  $\Delta C_p$  decreases with an increase of molecular weight and finally approaches to an asymptotic constant (for PS, the value,  $\Delta C_{p\infty}$  is 0.35 J/g K). The relationship between  $\Delta C_p$  and molecular weight of PS can be fitted by the following equation

$$\Delta C_p = [1 + (D_e/D_w)^{2/3}] \Delta C_{p\infty}, \quad (1)$$

where  $D_e$  is the number of the end-groups in a polymer molecule and  $D_w$  is the weight-average degree of polymerization. For linear polymer,  $D_e = 2$ ,  $D_w = M_w/M$ , where  $M$  is the molecular weight of a repeat unit and  $M_w$  is the weight-average molecular weight. In equation (1), it is obvious that  $\Delta C_p$  is related to the ratio of  $D_e$  and  $D_w$ , i.e. the larger the ratio, the larger the  $\Delta C_p$ . Therefore,  $\Delta C_p$  strongly depends on the concentration of the free end-groups in bulk polymer.

### Binary mixtures of DMWPS

For binary miscible mixtures of polymers, the dependence of  $T_g$  on composition can be presented by

three well-known equations [1, 2, 5]:

$$\text{Fox equation} \quad \frac{1}{T_g} = \frac{w_1}{T_{g1}} + \frac{w_2}{T_{g2}} \quad (2)$$

Couchman equation

$$\ln T_g = \frac{w_1 \ln T_{g1} + k w_2 \ln T_{g2}}{w_1 + k w_2} \quad (3)$$

Gordon-Taylor equation

$$T_g = \frac{w_1 T_{g1} + k w_2 T_{g2}}{w_1 + k w_2}, \quad (4)$$

where  $w_i$  ( $i = 1, 2$ ) is the weight fraction of component  $i$ ;  $T_{gi}$  ( $i = 1, 2$ ) is the glass transition temperature of component  $i$ ;  $T_g$  is the glass transition temperature of a mixture;  $k$  is equal to  $\Delta C_{p2}/\Delta C_{p1}$ . In these equations the interaction between components is not considered ( $\chi = 0$ ). Therefore, it is expected that these equations would give a good fit to results for mixtures having  $\chi = 0$ . However, this is far away from our expectation by comparing predictions with the experimental results. It is evident that the experimental curves in Figs 3 and 4 show a clear negative deviation from the calculated values using the above equations. Further investigation is in progress at our laboratory.

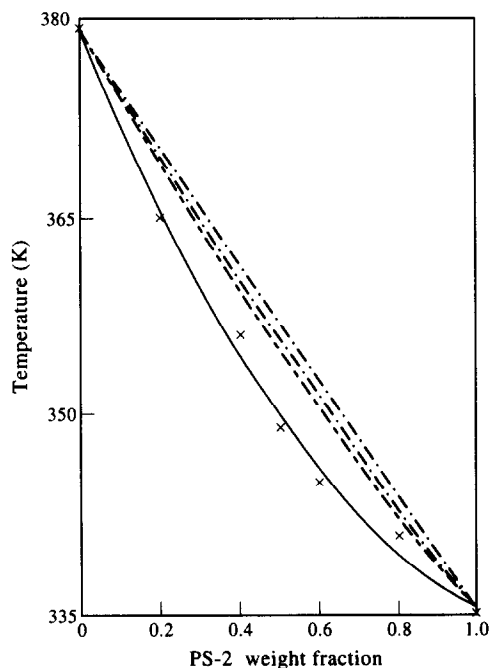


Fig. 4. Dependence of  $T_g$  on composition in PS-2/PS-11 mixtures: x experimental values; --- additive curve; — experimental curve; ··· Fox equation curve; -·- Gordon-Taylor equation curve for  $k = 1.13$ ; - - - Couchman equation curve for  $k = 1.13$ .

*Binary blends of PPO and DMWPS*

The blends PPO/PS are compatible at a segmental level [17]. The curves of  $T_g$  vs composition of the PPO/DMWPS blends are shown in Figs 5(a)–(e). According to observations, it can be known that the experimental values measured by DSC fit quite well with those calculated from the Couchman, Gordon–

Taylor and Fox equations with the increase of molecular weight of PS. When the molecular weight of PS is very low, the measured values are basically identical with the results of the Couchman equation. With an increase of molecular weight of PS, the experimental results gradually drift away from the values of the Couchman equation and fit those of the Gordon–Taylor equation. As the molecular weight of

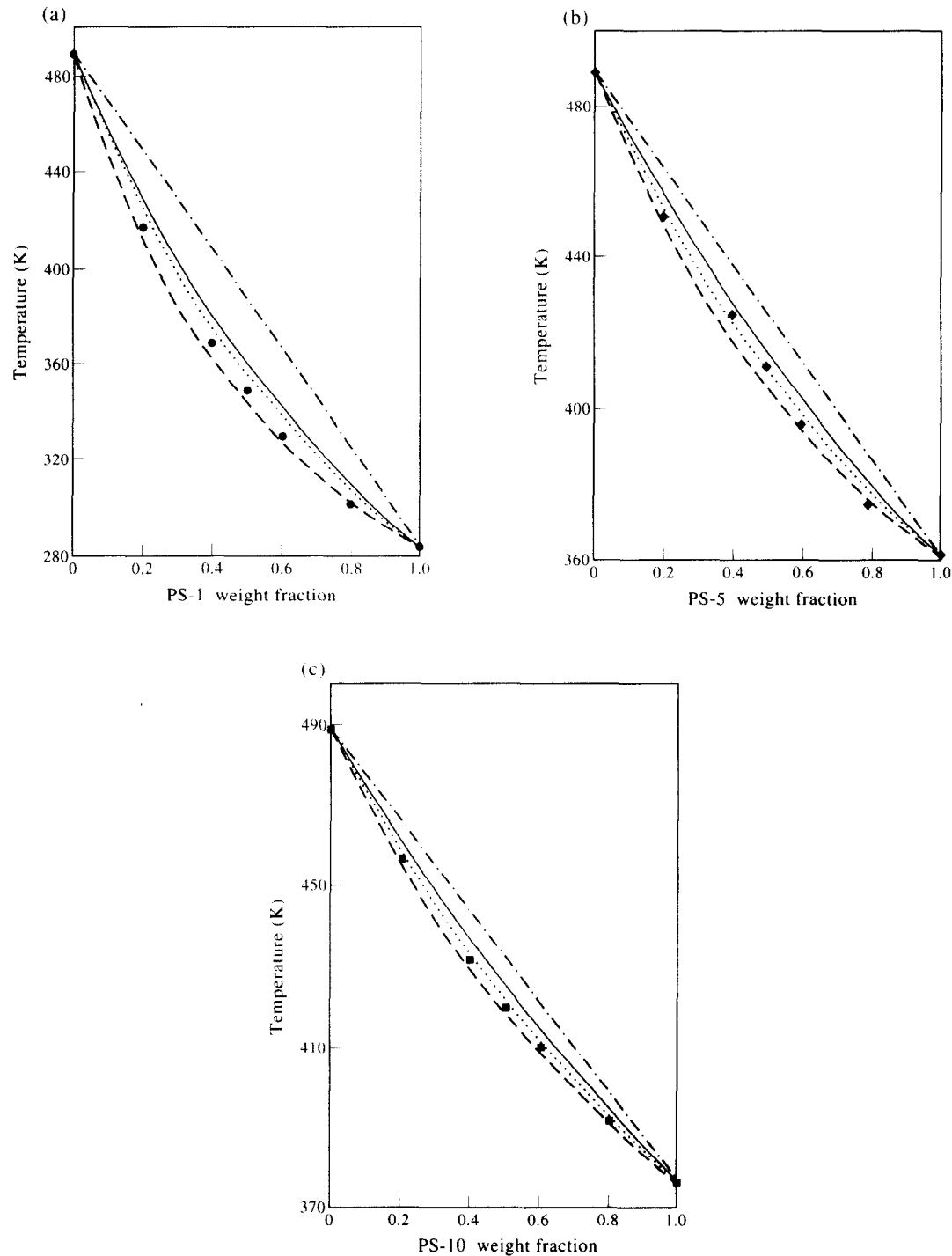


Fig. 5 (a–c) Caption on next page.

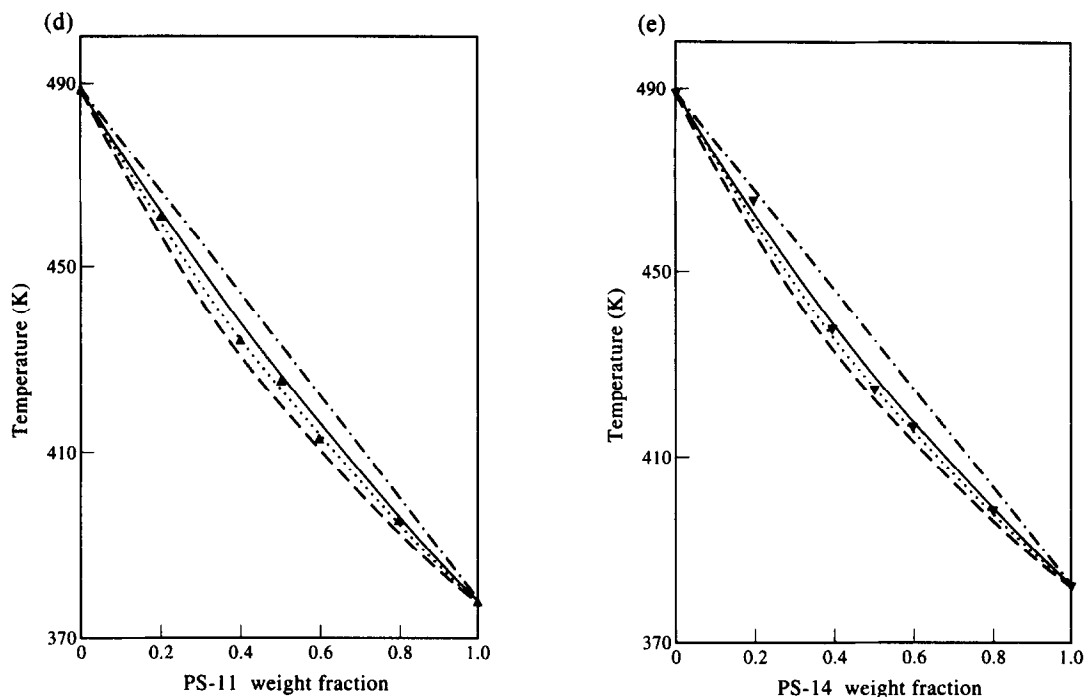


Fig. 5. Dependence of  $T_g$  on composition in PPO/DMWPS blends: (a) PPO/PS-1; (b) PPO/PS-5; (c) PPO/PS-10; (d) PPO/PS-11; (e) PPO/PS-14. ●, ◆, ■, ▲ and ▼ represent the experimental values of  $T_g$  in PPO/PS-1, PPO/PS-5, PPO/PS-10, PPO/PS-11 and PPO/PS-14, respectively. --- additive curve; — Fox equation curve; ··· Gordon-Taylor equation curve; -- Couchman equation curve.  $k$  is shown in Table 3.

PS is very high ( $M_w > 5.8 \times 10^5$ ), the Gordon-Taylor and Fox equations can give good fit to the experimental values. We think that the trend of  $T_g$  composition and the molecular weight of PS is related to the  $\Delta C_p$  of PS, because the parameter,  $k$ , is dependent on  $\Delta C_p$  of PS and PPO, where  $\Delta C_p$  is a constant. The Gordon-Taylor equation could cover a broader range. Although the blends PPO/PS have a negative interaction parameter and are different from the mixtures of PS ( $\chi = 0$ ), the calculated curves from equation (1), (2) and (3) are coincident with the experimental curves. Therefore, it is expected that the above equations are useful for polymer blends with a negative value of  $\chi$  (for PPO/PS blends,  $\chi = -0.06$  [18, 19]).

#### Effects of $\chi$ on $T_g$ of blends

Lu and Weiss [12] found that the Fox and Couchman equations, including the Gordon-Taylor equation, was not suitable to predict  $T_g$  for miscible systems with strong-specific interactions (for example, blends of polyamide-6 and manganese sulfonated polystyrene). They presented a modified equation for this system, i.e.

$$T_g = \frac{w_1 T_{g1} + w_2 T_{g2}}{w_1 + k w_2} - \frac{\chi R (T_{g2} - T_{g1}) b w_1 w_2}{\Delta C_{p1} (w_1 + k w_2) (w_1 + b w_2)^2}, \quad (5)$$

where  $b$  is the ratio of amorphous densities of polymer 2 and polymer 1;  $R$  is the gas constant. In equation (5),  $T_{g2} > T_{g1}$ . If  $T_{g2} < T_{g1}$ ,  $(T_{g2} - T_{g1})$  in equation (5) is replaced by  $(T_{g1} - T_{g2})$ , for PPO/PS blends, the specific interaction was found to be

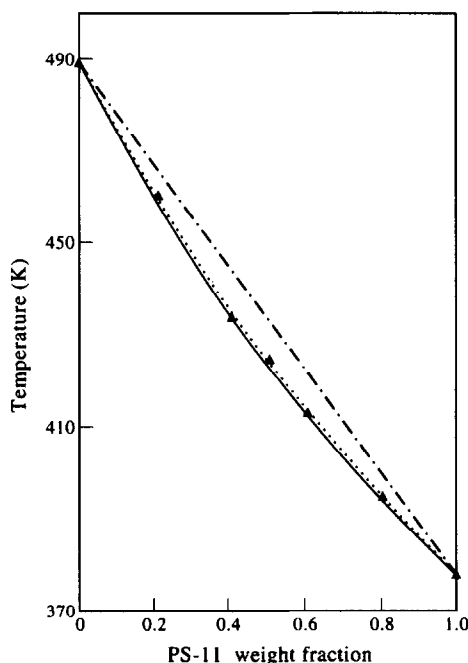


Fig. 6. Dependence of  $T_g$  on composition in PPO/PS-11 blends: ▲ experimental values; --- additive curve; — Gordon-Taylor equation curve; ··· Lu-Weiss equation curve.

$\chi = -0.06$  [18, 19]. Figure 6 gives the curves of the Gordon–Taylor and Lu–Weiss equations and the experimental values. It is shown that the calculated results of the Gordon–Taylor equation is basically identical to those of the Lu–Weiss equation. It results because the interaction between components is not so strong and so the increase of  $T_g$  is not evident. As seen from equation (5) itself, the maximum of values of  $\chi$ 's coefficient in the second term (including negative sign) is only about  $-7.4$  K. Since  $\chi$  is too small ( $= -0.06$ ), the maximum of  $T_g$ 's increment is only  $0.45$  K. Therefore, it is suggested that the Gordon–Taylor equation should be accurate enough to calculate  $T_g$  of blends with not very strong specific interactions and that the Lu–Weiss equation should be used for polymer blends with strong-specific interactions.

#### CONCLUSIONS

Our results suggest that the Fox, Couchman and Gordon–Taylor equations could be applied to polymer blends if the interaction between their molecules is not too strong. For a system with strong-specific interactions, the Lu–Weiss equation is a better formula. However, it is also proposed that the above equations could not be applied to mixtures when the interaction between components was very weak ( $\chi = 0$ ). It is also found that the Fox, Couchman and Gordon–Taylor equations are only suitable for describing the experimental results for a different molecular weight range of components.

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